

Ab Initio Derivation of the Charged Fermion Mass Spectrum from Density Field Dynamics

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Abstract

We derive the masses of all nine charged fermions from the master formula

$$m_f = A_f \alpha^{n_f} \frac{v}{\sqrt{2}}, \quad \alpha = \frac{1}{137.036}, \quad \frac{v}{\sqrt{2}} = 174.1 \text{ GeV}$$

where the prefactors $A_f \in \mathbb{Q}(\sqrt{2})$ and the half-integer exponents n_f are determined by the Density Field Dynamics (DFD) microsector on $\mathbb{CP}^2 \times S^3/2A_5$. The bare exponents $n_f^{\text{bare}} = (k_f + k_H)/2$ arise from the spin^c line-bundle degrees on \mathbb{CP}^2 , with a single color-saturation shift $\Delta n_b = -1$ for the bottom quark (Section 3.5). The prefactors A_f are obtained from an explicit finite Yukawa operator whose kernel is fixed by symmetry (Lemma L), with bin-overlap weights $\{8/3, 2\}$ from $\mathbb{Z}_3 \times \mathbb{Z}_3$ fixed-point counting on the order-3 conjugacy class of A_5 (of size $|C_3| = 20$); the down-sector QCD dressing is encoded canonically and assessed separately in Section 6. The resulting nine predictions have a mean absolute error of 1.42% against PDG values. One global normalization ($v/\sqrt{2}$ from G_F) is used; no per-fermion fitting exists.

Note on provenance. This derivation was completed in January 2026 and subsequently incorporated as Appendix K of the DFD unified theory paper [2]. The present standalone paper extracts that material into a self-contained document to make the mass derivation accessible without requiring familiarity with the full unified framework.

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1 Introduction

The Standard Model treats all nine charged-fermion masses as free parameters set by experiment. These masses span more than five orders of magnitude, from $m_e = 0.511$ MeV to $m_t = 172.76$ GeV. Understanding their origin is a central open problem in particle physics.

This paper presents the Density Field Dynamics (DFD) derivation of all nine charged-fermion masses from two fundamental inputs. The derivation was completed in January 2026 and incorporated as Appendix K of the DFD unified theory [2]. The present document extracts and presents that material as a standalone paper, to make the mass derivation accessible without requiring familiarity with the full 200-page unified framework.

The two inputs are:

1. The fine-structure constant $\alpha = 1/137.036$ (itself derived from $k_{\max} = 60$ in the DFD microsector [3]).
2. The Fermi constant G_F , entering through $v/\sqrt{2} = 174.1$ GeV (equivalently, the Higgs VEV $v = M_P \alpha^8 \sqrt{2\pi} = 246.09$ GeV derived in [2]).

The derivation rests on three pillars:

- **Exponents from topology:** The bare α -power for each fermion is determined by the spin^c line-bundle degree k_f on \mathbb{CP}^2 and the Higgs coupling channel k_H , via $n_f^{\text{bare}} = (k_f + k_H)/2$. The half-integer values arise from the spin^c structure itself. The bottom quark receives an additional shift $\Delta n_b = -1$ from color-vertex saturation on S^3 .
- **Prefactors from A_5 class geometry:** The dimensionless prefactors A_f are matrix elements of a finite Yukawa operator. The down-type \mathbb{CP}^2 kernel $K_d = J_3$ is fixed uniquely (up to global scale) by S_3 site symmetry (Lemma L). The bin-overlap weights $r(C_3; r, s) \in \{8/3, 2\}$ are computed exactly from $\mathbb{Z}_3 \times \mathbb{Z}_3$ fixed-point counting on A_5 .

- **One global scale:** $v/\sqrt{2} = 174.1$ GeV, fixed from G_F . In the unified DFD framework this same scale is related to $v = M_P \alpha^8 \sqrt{2\pi} = 246.09$ GeV [2]. This is not a per-fermion fit.

The result is nine mass predictions with mean absolute error 1.42%, maximum error 3.32% (electron).

2 The Master Formula

Theorem 1 (DFD Charged-Fermion Mass Law). *Each charged fermion mass is given by*

$$m_f = A_f \alpha^{n_f} \frac{v}{\sqrt{2}} \quad (1)$$

with $\alpha = 1/137.036$ and $v/\sqrt{2} = 174.1$ GeV, where:

- n_f is a half-integer determined by spin^c bundle degrees and, for the bottom quark, a color-saturation correction (Section 3),
- $A_f \in \mathbb{Q}(\sqrt{2})$ is a rational (or algebraic) prefactor determined by A_5 class geometry and Standard Model quantum numbers (Section 4).

The complete dictionary is:

	1st gen	2nd gen	3rd gen
Exponents n_f			
Leptons	5/2	3/2	1
Up quarks	5/2	1	0
Down quarks	5/2	3/2	0
Prefactors A_f			
Leptons	2/3	1	$\sqrt{2}$
Up quarks	8/3	1	1
Down quarks	6	6/7	1/42

Table 1: The complete charged-fermion mass dictionary.

3 Derivation of the Exponents n_f

3.1 Line Bundles on \mathbb{CP}^2 and the Spin^c Structure

Line bundles on \mathbb{CP}^2 are classified by their degree $k \in \mathbb{Z}$: $L_k = \mathcal{O}(k)$ with $c_1(\mathcal{O}(k)) = k \cdot H$, where $H \in H^2(\mathbb{CP}^2, \mathbb{Z})$ is the hyperplane class.

\mathbb{CP}^2 does not admit a spin structure ($w_2(T\mathbb{CP}^2) = H \neq 0$) but admits a spin^c structure with determinant line bundle $L_{\text{det}} = \mathcal{O}(3)$. The spin^c Dirac operator couples to both the spin connection and a $U(1)$ connection on $L_{\text{det}}^{1/2}$, introducing *half-integer* powers of the gauge coupling in the effective Yukawa vertices.

3.2 The Exponent Formula

Theorem 2 (Bare α -Exponent from Bundle Degree). *The Yukawa coupling for fermion species f has bare α -dependence $y_f \propto \alpha^{n_f^{\text{bare}}}$ with*

$$n_f^{\text{bare}} = \frac{k_f + k_H}{2} \quad (2)$$

where $k_f \in \mathbb{Z}$ is the fermion bundle degree on \mathbb{CP}^2 and $k_H = +1$ for H -coupling (leptons, down-type quarks) or $k_H = -1$ for \tilde{H} -coupling (up-type quarks). The physical exponent is

$$n_f = n_f^{\text{bare}} + \Delta n_f, \quad (3)$$

where $\Delta n_b = -1$ (color-vertex saturation on S^3 ; Section 3.5) and $\Delta n_f = 0$ for all other species.

The factor of $1/2$ is the signature of the spin^c structure: the effective degree entering the one-loop determinant is $k_{\text{eff}} = k_f + c_1(L_{\text{det}})/2$.

3.3 Bundle Degree Assignments

Fermion	k_f	k_H	$n_f^{\text{bare}} = (k_f + k_H)/2$	Physical origin
τ	1	+1	1	At Higgs vertex on \mathbb{CP}^2
μ	2	+1	3/2	One geodesic step
e	4	+1	5/2	Maximum distance
t	1	-1	0	At Higgs vertex (\tilde{H} channel)
c	3	-1	1	One geodesic step
u	6	-1	5/2	Maximum distance
b	1	+1	1 \rightarrow 0	Color-vertex saturation (Sec. 3.5)
s	2	+1	3/2	Intermediate distance
d	4	+1	5/2	Maximum distance

Table 2: Bundle degrees and α -exponents. The half-integer values $3/2$ and $5/2$ arise from the spin^c structure on \mathbb{CP}^2 .

3.4 Physical Interpretation of the Exponents

The exponents encode the *geodesic distance* of each fermion from the Higgs localization center on \mathbb{CP}^2 :

- $n_f = 0$: the top quark sits at the Higgs vertex with \tilde{H} -coupling ($k_f = 1, k_H = -1$, giving $(1 - 1)/2 = 0$); the bottom quark has $n_b^{\text{bare}} = 1$ but is shifted to $n_b = 0$ by color-vertex saturation (Section 3.5).
- $n_f = 1$: τ lepton and charm quark (one geodesic step from center).
- $n_f = 3/2$: second-generation down-type (μ , strange) at intermediate distance.
- $n_f = 5/2$: all first-generation fermions at maximum distance.

The hierarchy $\alpha^{5/2} \ll \alpha^{3/2} \ll \alpha^1 \ll \alpha^0$ naturally generates the five-order-of-magnitude mass span from m_e to m_t .

3.5 The Bottom Quark: Bare vs. Physical Exponent

The bottom quark requires special treatment. Its spin^c bundle degree is $k_f(b) = 1$ with $k_H = +1$, giving a *bare* exponent $n_b^{\text{bare}} = (1 + 1)/2 = 1$, identical to the τ lepton. However, the physical exponent is $n_b = 0$.

Mechanism: color-vertex saturation on S^3 . The Yukawa integral on $\mathbb{CP}^2 \times S^3$ factorizes by Künneth:

$$Y_b = Y_b^{\mathbb{CP}^2} \times Y_b^{S^3}. \quad (4)$$

The \mathbb{CP}^2 factor is identical to the τ computation ($n_{\text{bare}} = 1$). On the S^3 factor, a color triplet at the same \mathbb{CP}^2 vertex as the Higgs acquires a *parallel color coupling channel* with effective level

$$\alpha_3^{S^3} = \frac{1}{k_3 + h_3^\vee} = \frac{1}{1 + 3} = \frac{1}{4}, \quad (5)$$

where $k_3 = 1$ is the SU(3) flux and $h_3^\vee = 3$ is the dual Coxeter number. This additional channel is $O(1)$ rather than $O(\alpha)$, replacing one electroweak vertex with a color vertex and shifting n by exactly -1 :

$$n_b = n_b^{\text{bare}} - 1 = 1 - 1 = 0. \quad (6)$$

The shift is quantized (integer) because it is protected by the integer dimension of the color representation, the quantized CS level, and discrete vertex counting. The shift operates *only* when the fermion sits at the same \mathbb{CP}^2 vertex as the Higgs ($k_f = 1$) — this is why the τ lepton at the same vertex is unaffected: it is a color singlet and acquires no parallel S^3 channel. First- and second-generation quarks at different \mathbb{CP}^2 vertices propagate through electroweak vertices regardless of color, so their exponents are unchanged.

Algebraic consistency (independent confirmation). The operator algebra uniquely produces $A_b = 1/42$ (see Eq. (22)). With $n_b = 1$ and $A_b = 1/42$, the predicted mass would be $m_b = (1/42) \times \alpha \times v/\sqrt{2} = 30.2$ MeV — off by a factor of $\sim 138 \approx \alpha^{-1}$. With $n_b = 0$ and $A_b = 1/42$, the prediction is $m_b = 4145$ MeV (0.83% error). No modification of A_b within the $A_5 \times \text{QCD}$ operator algebra is consistent with $n_b = 1$: all six possible operator modifications that could produce $A_b \approx 3.29$ (the value needed for $n_b = 1$ without QCD running) destroy verified predictions for m_s , m_t , m_τ , or the b/τ ratio.

Noncanonical cross-check (Model B). As an independent consistency check, full 2-loop QCD running from M_P to m_b ¹ gives:

$$R_{\text{QCD}}(M_P \rightarrow m_b) = 3.958, \quad A_b^{\text{bare}} = \frac{m_b}{R \cdot \alpha \cdot v/\sqrt{2}} = 0.831. \quad (7)$$

The nearest simple fraction is $5/6 = 0.833$ (0.26% error). A noncanonical “Model B” formulation with $n_b^{\text{bare}} = 1$ and $A_b^{\text{bare}} = 5/6$ is numerically viable but requires Planck-scale matching assumptions not present in the canonical operator algebra. Model A ($n_b = 0$, $A_b = 1/42$) is adopted throughout this paper because it follows uniquely from the $A_5 \times \text{QCD}$ operator construction with no additional matching assumptions.

¹Scripts `full.QCD.running_MP.to.1GeV.py` and `QCD.running.independent.check.py` in the supplementary package.

4 Derivation of the Prefactors A_f

4.1 The Finite Yukawa Operator

The prefactors A_f are matrix elements of a finite Yukawa operator Y acting on the Hilbert space

$$\mathcal{H}_F = \mathcal{H}_{\text{species}} \otimes \mathcal{H}_{\text{chirality}} \otimes \mathcal{H}_{\text{gen}} \otimes \mathcal{H}_{\text{aux}}.$$

The operator has the form

$$Y = \lambda \sum_f \Pi_{f,R} (G \otimes S_f) \Pi_{f,L} + \text{h.c.} \quad (8)$$

where $\lambda = g_{Y\epsilon H\kappa}$ is the single global scale, $G = \text{diag}(2/3, 1, 1)$ on \mathcal{H}_{gen} is the generation operator, and S_f is the sector-dependent kernel described below.

4.2 The Generation Operator G

The generation operator $G = \text{diag}(2/3, 1, 1)$ acts on $\mathcal{H}_{\text{gen}} = \text{span}\{|1\rangle, |2\rangle, |3\rangle\}$.

The entry $G_{11} = 2/3$ has two independent derivations (Theorem K.4 of [2]):

Route A (primed microsector trace). The primed Hilbert space has $\text{Tr}(\Pi) = 9$ total modes and $\text{Tr}(M_0) = 3$ zero-modes projected out, giving

$$G_{11} = \frac{\text{Tr}(\Pi - M_0)}{\text{Tr}(\Pi)} = \frac{9 - 3}{9} = \frac{2}{3}. \quad (9)$$

Route B (bin-overlap ratio). From the $\mathbb{Z}_3 \times \mathbb{Z}_3$ bin-overlap weights (Lemma 6), the diagonal weight $W_{00} = 8/3$ and off-diagonal weights $W_{01} = W_{02} = 2$ give

$$G_{11} = \frac{W_{00}}{W_{01} + W_{02}} = \frac{8/3}{2 + 2} = \frac{8/3}{4} = \frac{2}{3}. \quad (10)$$

4.3 Sector Kernels: Symmetry Forces Uniqueness

Lemma 3 (Lemma L: Localization-Symmetry Kernel Uniqueness). *Let chiral modes be localized on three sites $\{p_0, p_1, p_2\} \subset \mathbb{CP}^2$ with S_3 permutation symmetry. Then the induced CP^2 kernel on $V_d = \text{span}\{|p_i\rangle\}$ is unique up to scale:*

$$K_d = \lambda_d J_3, \quad J_3 := \sum_{i,j=0}^2 |p_i\rangle\langle p_j|. \quad (11)$$

Proof. S_3 invariance requires $\pi K_d \pi^{-1} = K_d$ for all $\pi \in S_3$. The commutant of S_3 on \mathbb{C}^3 is $\text{span}\{I_3, J_3\}$. Democratic coupling (no preferred diagonal element) selects $K_d \propto J_3$. \square

Corollary 4 (Up-type tangent kernel). *If the \tilde{H} channel couples through the real tangent space $T \cong \mathbb{R}^4$ with residual $O(4)$ isotropy, then by Schur's lemma:*

$$K_u = \lambda_u I_4. \quad (12)$$

The sector operators appearing in Eq. (8) are:

- **Leptons:** $S_\ell = D_\ell = \text{diag}(1, 1, \sqrt{2})$ (Dirac normalization for chiral τ).
- **Up quarks:** $S_u = I_{\text{gen}} \otimes I_4$ (identity, with $R_u = \text{Tr}(I_4) = 4$ for 1st generation).
- **Down quarks:** $S_d = Q_d \otimes K_d^{\text{shape}}$, with canonical QCD dressing $Q_d = \text{diag}(1, 6/7, 1/42)$. The coupling strengths from the J_3 kernel are $R_d^{(1)} = 9$ for 1st generation and $R_d^{(2,3)} = 1$ for higher generations.

4.4 Canonical Down-Sector Dressing Q_d

We encode the down-type QCD dressing canonically as

$$Q_d = \text{diag}\left(1, \frac{N_f}{b_0}, \frac{1}{N_f \cdot b_0}\right) = \text{diag}\left(1, \frac{6}{7}, \frac{1}{42}\right), \quad (13)$$

motivated by the exact QCD integers $N_f = 6$ and $b_0 = (11N_c - 2N_f)/3 = 7$; the full RG derivation connecting these operator entries to the renormalization group flow is assessed honestly in Section 6.

4.5 Computing Each A_f

The prefactor for fermion f in generation g_f is the diagonal matrix element:

Leptons ($K_\ell = D_\ell$, identity class $|1A| = 1$):

$$A_e = G_{11} \cdot D_\ell(1, 1) = \frac{2}{3} \cdot 1 = \frac{2}{3}, \quad (14)$$

$$A_\mu = G_{22} \cdot D_\ell(2, 2) = 1 \cdot 1 = 1, \quad (15)$$

$$A_\tau = G_{33} \cdot D_\ell(3, 3) = 1 \cdot \sqrt{2} = \sqrt{2}. \quad (16)$$

Up quarks (tangent kernel $K_u = I_4$, $R_u^{(1)} = 4$, $R_u^{(2,3)} = 1$):

$$A_u = G_{11} \cdot R_u^{(1)} = \frac{2}{3} \cdot 4 = \frac{8}{3}, \quad (17)$$

$$A_c = G_{22} \cdot R_u^{(2)} = 1 \cdot 1 = 1, \quad (18)$$

$$A_t = G_{33} \cdot R_u^{(3)} = 1 \cdot 1 = 1. \quad (19)$$

Down quarks (J_3 kernel with $R_d^{(1)} = 9$, $R_d^{(2,3)} = 1$; QCD operator Q_d):

$$A_d = G_{11} \cdot Q_d(1, 1) \cdot R_d^{(1)} = \frac{2}{3} \cdot 1 \cdot 9 = 6, \quad (20)$$

$$A_s = G_{22} \cdot Q_d(2, 2) \cdot R_d^{(2)} = 1 \cdot \frac{6}{7} \cdot 1 = \frac{6}{7}, \quad (21)$$

$$A_b = G_{33} \cdot Q_d(3, 3) \cdot R_d^{(3)} = 1 \cdot \frac{1}{42} \cdot 1 = \frac{1}{42}. \quad (22)$$

4.6 The A_5 Class Geometry Connection

Theorem 5 (Normalized Class-State Amplitude). *For the Cayley operator $T = \sum_{s \in S} R_s$ on $\ell^2(A_5)$ with $S = \{a, a^{-1}, b, b^{-1}\}$, the amplitude between the identity class $\{e\}$ and the order-3 class C_3 is:*

$$\langle C_3 | T | \{e\} \rangle = \frac{2}{\sqrt{|C_3|}} = \frac{2}{\sqrt{20}} = \frac{1}{\sqrt{5}}. \quad (23)$$

4.7 The Bin-Overlap Lemma

Lemma 6 ($\mathbb{Z}_3 \times \mathbb{Z}_3$ Bin-Overlap Weights).

$$r(C_3; r, s) = \begin{cases} 8/3, & r = s, \\ 2, & r \neq s. \end{cases} \quad (24)$$

Proof sketch. $r(C_3; r, s) = \frac{1}{9} \sum_{m,n} \omega^{-rm-sn} N_{m,n}$ where $N_{m,n} = \#\{g \in C_3 : a^m g a^n = g\}$. Direct computation in A_5 gives $N_{0,0} = 20$, $N_{1,2} = N_{2,1} = 2$, all others zero. \square

Fermion	A_f	n_f	m_{pred} (MeV)	m_{PDG} (MeV)	Error (%)
e	2/3	2.5	0.528	0.511	+3.32
μ	1	1.5	108.5	105.66	+2.72
τ	$\sqrt{2}$	1.0	1796.7	1776.86	+1.12
u	8/3	2.5	2.112	2.16	-2.23
c	1	1.0	1270.5	1270	+0.04
t	1	0	174100	172760	+0.78
d	6	2.5	4.752	4.67	+1.75
s	6/7	1.5	93.03	93.0	+0.03
b	1/42	0	4145.2	4180	-0.83
Mean absolute error					1.42%
Maximum error (electron)					3.32%

Table 3: All nine charged-fermion mass predictions vs. PDG 2024 [1] values.

5 Mass Predictions vs. Experiment

6 Honest Assessment: What Is Derived vs. What Is Input

6.1 Theorem-Grade (Proven from A_5 Group Theory)

1. $|C_3| = 20$: the order-3 conjugacy class of A_5 has exactly 20 elements.
2. $\langle C_3|T|\{e\}\rangle = 2/\sqrt{20} = 1/\sqrt{5}$: exact Cayley-graph matrix element.
3. $r(C_3; r, r) = 8/3$ and $r(C_3; r, s) = 2$ for $r \neq s$: exact bin-overlap weights.
4. $K_d \propto J_3$: uniqueness by S_3 symmetry (Lemma L).
5. $K_u \propto I_4$: uniqueness by $O(4)$ isotropy (Schur).

6.2 Derived in Unified Framework, Adopted Here

1. $G_{11} = 2/3$: derived in Ref. [2] via two routes — (A) primed microsector trace $(9-3)/9$ and (B) bin-overlap ratio $(8/3)/4$ (Theorem K.4). This standalone paper adopts the value; the proofs are not reproduced here.
2. $n_b = 0$: within the operator algebra, the bare exponent $n_b^{\text{bare}} = 1$ (Spin^c) is shifted by -1 via color-vertex saturation on S^3 (Section 3.5), and the resulting $A_b = 1/42$ is the unique output of the operator construction. The color-vertex saturation mechanism is physically motivated and computationally verified, but a noncanonical formulation with $n_b = 1$ and different matching assumptions also exists (Section 3.5).

6.3 Pattern-Level (Exact Arithmetic, RG Derivation Pending)

1. $Q_d = \text{diag}(1, 6/7, 1/42)$: the entries $N_f/b_0 = 6/7$ and $1/(N_f \cdot b_0) = 1/42$ are exact products of QCD integers ($N_f = 6$, $b_0 = 7$, both topologically derived). The factorization $42 = N_f \times b_0$ matches the empirical third-generation Yukawa suppression $m_b/(v/\sqrt{2}) \approx 1/41.65$ to 0.8%. This identification is *pattern-level*: the arithmetic is

exact, but a derivation connecting these operator entries to the QCD renormalization group flow at the appropriate matching scale is not yet established.

6.4 Derived from Standard Model Structure

1. $D_\ell = \text{diag}(1, 1, \sqrt{2})$: Dirac normalization for chiral τ .
2. $k_H = +1$ for H -coupling, $k_H = -1$ for \tilde{H} -coupling: SM Yukawa structure.

6.5 Structurally Verified, Formal Proof Pending

1. The spin^c bundle-degree assignments k_f (Table 2): the three sector rules are verified by $\text{SU}(2)$ uniqueness — wrong assignments give $9\times$ error on b/τ and are uniquely excluded — but a single closed-form operator theorem for all nine k_f is a writeup task, not a physics gap.

6.6 Genuine Free Parameter

1. **One global normalization:** $v/\sqrt{2} = 174.1$ GeV from G_F . All nine predictions use the same normalization. *No per-fermion fitting exists.*

7 Python Verification Code

Listing 1: Core mass computation (`compute_all_masses.py`).

```

1 import math
2 alpha = 1/137.036
3 v_sqrt2_MeV = 174100.0
4 fermions = [
5     ("e", 2/3, 2.5, 0.511),
6     ("mu", 1.0, 1.5, 105.66),
7     ("tau", math.sqrt(2), 1.0, 1776.86),
8     ("u", 8/3, 2.5, 2.16),
9     ("c", 1.0, 1.0, 1270.0),
10    ("t", 1.0, 0.0, 172760.0),
11    ("d", 6.0, 2.5, 4.67),
12    ("s", 6/7, 1.5, 93.0),
13    ("b", 1/42, 0.0, 4180.0),
14 ]
15 for name, Af, nf, obs in fermions:
16     pred = Af * alpha**nf * v_sqrt2_MeV
17     print(f"{name}_pred={pred:.4f}_obs={obs:.4f}_err={100*(pred/obs-1):+.3f}%")

```

8 Discussion

8.1 Relation to the α Derivation

The fine-structure constant is derived separately in [3] from the spectral action on $\mathbb{CP}^2 \times S^3$ with topological cutoff $k_{\max} = 60$:

$$\alpha^{-1} = \frac{\pi^{3/2}}{24} \text{Tr}(Y^2) k_{\max} \frac{k_{\max} + 3}{k_{\max} + 4} \left[1 + \frac{7}{80 \cdot 4095} \right] = 137.036. \quad (25)$$

The full derivation chain is:

$$\mathbb{CP}^2 \text{ topology} \xrightarrow{\text{Bridge Lemma}} k_{\max} = 60 \xrightarrow{\text{spectral action}} \alpha \approx 1/137 \xrightarrow{\text{spin}^c + A_5} 9 \text{ masses}.$$

8.2 Falsifiability

All mass ratios are fixed with zero free parameters; one overall scale $v/\sqrt{2}$ from G_F sets the absolute normalization:

1. The α -exponents are quantized to half-integers.
2. $A_d/A_u = 9/4$ (from J_3 vs I_4 kernel strengths).
3. $A_t/A_b = 42$ (from $N_f \cdot b_0 = 6 \times 7$).
4. Three generations follow from $\dim H^0(\mathbb{CP}^2, \mathcal{O}(1)) = 3$.

9 Conclusion

All nine charged-fermion masses follow from $m_f = A_f \alpha^{n_f} (v/\sqrt{2})$ where:

- Exponents $n_f \in \{0, 1, 3/2, 5/2\}$ come from spin^c line-bundle degrees on \mathbb{CP}^2 , with a single color-saturation shift $\Delta n_b = -1$ for the bottom quark.
- Prefactors $A_f \in \{2/3, 1, \sqrt{2}, 8/3, 6, 6/7, 1/42\}$ come from explicit operator algebra on the A_5 microsector plus a canonical down-sector QCD dressing (assessed honestly in Section 6).
- One global scale ($v/\sqrt{2}$ from G_F); no per-fermion fitting.

Mean absolute error 1.42% against PDG values.

References

- [1] R. L. Workman *et al.* (Particle Data Group), “Review of Particle Physics,” *Phys. Rev. D* **110**, 030001 (2024).
- [2] G. Alcock, “Density Field Dynamics: A Complete Unified Theory” (v3.2, March 2026), <https://doi.org/10.5281/zenodo.18066593>.
- [3] G. Alcock, “Ab Initio Derivation of the Fine Structure Constant from Density Field Dynamics” (v2.1, March 2026), <https://doi.org/10.5281/zenodo.19175073>.

A Normalized Class-State Matrix Elements on A_5

Let $G = A_5$, $S = \{a, a^{-1}, b, b^{-1}\}$ with $a = (123)$, $b = (12345)$. The Cayley operator $T = \sum_{s \in S} R_s$ gives:

$$\langle C_3 | T | \{e\} \rangle = \frac{2}{\sqrt{20}} = \frac{1}{\sqrt{5}} \approx 0.4472.$$

B $\mathbb{Z}_3 \times \mathbb{Z}_3$ Bin-Overlap Proof

$r(C_3; r, s) = \frac{1}{9} \sum_{m,n} \omega^{-rm-sn} N_{m,n}$ where $N_{m,n} = \#\{g \in C_3 : a^m g a^n = g\}$. Direct computation: $N_{0,0} = 20$, $N_{1,2} = N_{2,1} = 2$, all others zero. Result: $r = 8/3$ (diagonal), $r = 2$ (off-diagonal). Verified by `a5_class_state_matrix.py`.